

Two-body density operator $\hat{\rho}_2(\underline{x}_1, \underline{x}_2)$:

$$\begin{aligned}\hat{\rho}_2(\underline{x}_1, \underline{x}_2) \psi(\underline{R}) &= \left(\sum_{i=1}^N \sum_{\substack{j=1 \\ j \neq i}}^N \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) \psi(\underline{R}) \right) \\ &= \sum_{i \neq j=1}^N \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) \psi(\underline{R}) \quad (37)\end{aligned}$$

Two-body density

$$\langle \hat{\rho}_2(\underline{x}_1, \underline{x}_2) \rangle = \text{Tr} \{ \hat{\rho}_2(\underline{x}_1, \underline{x}_2) \hat{W} \} = \rho_2(\underline{x}_1, \underline{x}_2) \quad (38)$$

From equation (24);

$$\begin{aligned}\hat{\rho}_2(\underline{x}_1, \underline{x}_2) &= \int_{V^N} [\hat{\rho}_2(\underline{x}_1, \underline{x}_2) W(\underline{R}, \underline{R}')] (\underline{R}' = \underline{R}) d\underline{R} \\ &= \int_{V^N} \left[\sum_{i \neq j=1}^N \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) W(\underline{R}, \underline{R}') \right] (\underline{R}' = \underline{R}) d\underline{R} \\ &= \sum_{i \neq j=1}^N \int_{V^N} \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) \underbrace{W(\underline{R}, \underline{R})}_{=P(\underline{R})} d\underline{R} \\ &= \sum_{i \neq j=1}^N \int_{V^N} \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) P(\underline{r}_1, \underline{r}_2, \dots, \underline{r}_i, \dots, \underline{r}_j, \dots, \underline{r}_N) d\underline{r}_1, d\underline{r}_2, \dots, d\underline{r}_i, \dots, d\underline{r}_j, \dots, d\underline{r}_N \\ &= \sum_{i \neq j=1}^N \int_{V^N} \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) P(\underline{r}_i, \underline{r}_2, \dots, \underline{r}_1, \dots, \underline{r}_j, \dots, \underline{r}_N) d\underline{r}_i, d\underline{r}_2, \dots, d\underline{r}_1, \dots, d\underline{r}_j, \dots, d\underline{r}_N \\ &= \sum_{i \neq j=1}^N \int_{V^{N-2}} P(\underline{x}_1, \underline{x}_2, \dots, \underline{r}_1, \dots, \underline{r}_2, \dots, \underline{r}_N) d\underline{r}_3, \dots, d\underline{r}_1, \dots, d\underline{r}_2, \dots, d\underline{r}_N \\ &= \sum_{i \neq j=1}^N \int_{V^{N-2}} P(\underline{x}_1, \underline{x}_2, \underline{r}_3, \dots, \underline{r}_N) d\underline{r}_3, \dots, d\underline{r}_N\end{aligned}$$

In this last step, we are relabelling the integration variables $\underline{r}_1 = \underline{r}_i$, $\underline{r}_2 = \underline{r}_j$.

We get the same result for all $i \neq j$. It equals the number of ordered pairs

$$\underbrace{(i, j)}_{i \neq j} = N(N-1).$$

$$\begin{aligned}\rho_2(\underline{x}_1, \underline{x}_2) &= N(N-1) \int_{V^{N-2}} P(\underline{x}_1, \underline{x}_2, \underline{r}_3, \dots, \underline{r}_N) d\underline{r}_3 \dots d\underline{r}_N \\ &= N(N-1) \int_{V^{N-2}} W(\underline{x}_1, \underline{x}_2, \underline{r}_3, \dots, \underline{r}_N, \underline{x}_1, \underline{x}_2, \underline{r}_3, \dots, \underline{r}_N) d\underline{r}_3 \dots d\underline{r}_N \\ &= \left\langle \sum_{i \neq j=1}^N \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) \right\rangle \quad (39) \\ &\int_V \rho(\underline{x}) d\underline{x} = N \quad (35)\end{aligned}$$

The above can be extended for the two-body density $\rho_2(\underline{x}_1, \underline{x}_2)$;

$$\int_{V^2} \rho_2(\underline{x}_1, \underline{x}_2) d\underline{x}_i d\underline{x}_j = N(N-1) \quad (40)$$

because we have $N(N-1)$ distinct ordered pairs (i, j) with $1 \leq i \neq j \leq N$.

Reversion Relations

$$\int_V \rho_2(\underline{x}_1, \underline{x}_2) d\underline{x}_2 = (N-1)\rho(\underline{x}_1) \quad (41)$$

where N is the total number of particles.

$$\int_V \rho_2(\underline{x}_1, \underline{x}_2) d\underline{x}_1 = (N-1)\rho(\underline{x}_2) \quad (42)$$

Exchange Symmetry of Two-Body Density

$$\rho_2(\underline{x}_1, \underline{x}_2) = \rho_2(\underline{x}_2, \underline{x}_1) \quad (43)$$

$\rho_2(\underline{x}_1, \underline{x}_2)$ is the probability density for a particle to be at \underline{x}_1 and another at \underline{x}_2 .

$\rho_2(\underline{x}_1, \underline{x}_2) d\underline{x}_1 d\underline{x}_2$ is the probability for a particle to be in volume element

$d\underline{x}_1 = dx_1 dy_1 dz_1$ centered at \underline{x}_1 and another particle to be in volume element

$d\underline{x}_2 = dx_2 dy_2 dz_2$ centered on \underline{x}_2 .

Translational Invariance

For a liquid, translational invariance yields $\rho(\underline{x} + \underline{a}) = \rho(\underline{x})$ for any 3D translation \underline{a} .

$$\rightarrow \rho(\underline{x}) = \rho_0 = \text{const.}$$

$$\rho_2(\underline{x}_1 + \underline{a}, \underline{x}_2 + \underline{a}) = \rho_2(\underline{x}_1, \underline{x}_2)$$

for any 3D translation \underline{a} .

Take $\underline{a} = -\underline{x}_2$;

$$\rho_2(\underline{x}_1, \underline{x}_2) = \rho_2(\underline{x}_1 - \underline{x}_2, 0) = \rho_2(\underline{x}_1 - \underline{x}_2)$$

For a liquid, $\rho_2(\underline{x}_1, \underline{x}_2)$ is a function of the relative coordinates $\underline{r} = \underline{x}_1 - \underline{x}_2$ only.

$$\rho_2(\underline{x}_1, \underline{x}_2) = \rho_2(\underline{x}_1 - \underline{x}_2) = \rho_2(\underline{r}) \quad (44)$$

where $\underline{r} = \underline{x}_1 - \underline{x}_2$.

Potential Energy for a Liquid

From equation (25);

$$\begin{aligned} E_{pot} &= \int_{V^N} V(\underline{R}) W(\underline{R}, \underline{R}) d\underline{R} \\ &= \frac{1}{2} \sum_{i \neq j=1}^N \int_{V^N} v(\underline{r}_{ij}) W(\underline{R}, \underline{R}) d\underline{R} \end{aligned}$$

where $v(\underline{r}_{ij}) = v(|\underline{r}_i - \underline{r}_j|)$.

$$\begin{aligned} E_{pot} &= \frac{1}{2} \int_{V^2} v(|\underline{x}_1 - \underline{x}_2|) \left[\underbrace{\int_{V^N} \left\{ \sum_{i \neq j=1}^N \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) \right\} W(\underline{R}, \underline{R}) d\underline{R}}_{= \left\langle \sum_{i \neq j=1}^N \delta(\underline{x}_1 - \underline{r}_i) \delta(\underline{x}_2 - \underline{r}_j) \right\rangle} \right] d\underline{x}_1 d\underline{x}_2 \\ &= \left\langle \rho_2(\underline{r}_1, \underline{r}_2) \right\rangle \\ &= \rho_2(\underline{x}_1, \underline{x}_2) \end{aligned}$$

$$= \frac{1}{2} \int_{V^2} \rho_2(\underline{x}_1, \underline{x}_2) v(|\underline{x}_1 - \underline{x}_2|) d\underline{x}_1 d\underline{x}_2 \quad (45)$$

For a liquid,

$$E_{pot} = \frac{1}{2} \int_{V^2} \rho_2(\underline{x}_1 - \underline{x}_2) v(|\underline{x}_1 - \underline{x}_2|) d\underline{x}_1 d\underline{x}_2$$

$$\rho_2(\underline{x}_1, \underline{x}_2) = \rho_2(\underline{x}_1 - \underline{x}_2) = \rho_2(|\underline{r}|) = \rho_2(|\underline{x}_1 - \underline{x}_2|)$$

Rotational invariance of a uniform liquid – no directional dependence.

$$\rho_0 = \frac{N}{V} = \rho(\underline{x}) = \text{const.}$$

where ρ_0 is the bulk particle number density.

$$\rho_2(r) = \rho_0^2 g(r)$$

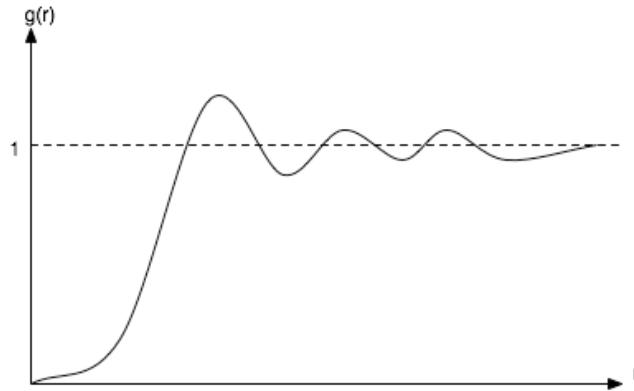
where $g(r)$ is the radial distribution function,

$$g(r) = \frac{\rho_2(r)}{\rho_0^2}$$

Let $\epsilon_{pot} = \frac{E_{pot}}{N}$ be the potential energy per particle. Then,

$$\epsilon_{pot} = 2\pi\rho_0 \int_0^\infty g(r)v(r)dr \quad (46)$$

where the 2π is due to the factor $\frac{1}{2}$ in E_{pot} .



$$\xrightarrow[\lim_{|\underline{x}_1 - \underline{x}_2| \rightarrow \infty}]{} \rho_2(\underline{x}_1, \underline{x}_2) = \rho(\underline{x}_1)\rho(\underline{x}_2) = \rho_0^2$$

as particles don't interact over a large distance. The last part of this is true for a liquid.

$$\xrightarrow[r \rightarrow \infty]{} \frac{\rho_2(r)}{\rho_0^2} = \xrightarrow[r \rightarrow \infty]{} \frac{\rho_0^2}{\rho_0^2} = 1 = \xrightarrow[r \rightarrow \infty]{} g(r)$$

3. Path Integral Representation of the Density Matrix

Coordinate Space Representation of Operators = Coordinate Space Elements of Operators (more generally).

$$\hat{O}_{\underline{R}}\psi(\underline{R}) = \tilde{\psi}(\underline{R}) = \int_{V^N} \tilde{\psi}(\underline{R}') \delta(\underline{R}' - \underline{R}) d\underline{R}'$$

where $\delta(\underline{R}' - \underline{R})$ is a 3N-dimensional delta function.

$$\begin{aligned} \tilde{\psi}(\underline{R}') &= \hat{O}_{\underline{R}}\psi(\underline{R}') \\ &= \int_{V^N} [\hat{O}_{\underline{R}}\psi(\underline{R}')] \delta(\underline{R}' - \underline{R}) d\underline{R}' \\ &= \int_{V^N} \psi(\underline{R}') [\hat{O}_{\underline{R}}\delta(\underline{R}' - \underline{R})] d\underline{R}' \end{aligned}$$

$$O(\underline{R}, \underline{R}') = \hat{O}_{\underline{R}} \delta(\underline{R}' - \underline{R})$$

(Apply operator \hat{O} to variable \underline{R}' in the three-dimensional delta function $\delta(\underline{R}' - \underline{R})$)

$$\hat{O}\psi(\underline{R}) = \int_{V^N} O(\underline{R}, \underline{R}') \psi(\underline{R}') d\underline{R}' \quad (47)$$

$$O(\underline{R}, \underline{R}') = \hat{O}_{\underline{R}} \delta(\underline{R}' - \underline{R})$$

$O(\underline{R}, \underline{R}')$ is the matrix elements of \hat{O} in coordinate space representation of \hat{O} .

$$\begin{aligned} \delta(\underline{R}' - \underline{R}) &= \delta(\underline{r}_1' - \underline{r}_1) \delta(\underline{r}_2' - \underline{r}_2) \dots \delta(\underline{r}_N' - \underline{r}_N) \\ &= \delta(x_1' - x_1) \delta(y_1' - y_1) \delta(z_1' - z_1) \delta(x_2' - x_2) \delta(y_2' - y_2) \delta(z_2' - z_2) \\ &\dots \delta(x_N' - x_N) \delta(y_N' - y_N) \delta(z_N' - z_N) \end{aligned}$$

This is the $3N$ -dimensional Dirac delta function.

$$\begin{aligned} \int_{V^N} \delta(\underline{R}' - \underline{R}) \psi(\underline{R}') d\underline{R}' &= \psi(\underline{R}) \\ &= \int_{V^N} \prod_{i=1}^N \delta(\underline{r}_i' - \underline{r}_i) \psi(\underline{r}_1', \underline{r}_2', \dots, \underline{r}_N') d\underline{r}_1' d\underline{r}_2' \dots d\underline{r}_N' \\ &= \psi(\underline{r}_1, \underline{r}_2, \dots, \underline{r}_N) \quad (49) \end{aligned}$$

$$\int_{V^N} \delta(\underline{R}' - \underline{R}) \psi(\underline{R}) d\underline{R} = \psi(\underline{R}')$$

$$\int_{V^N} \delta(\underline{R}' - \underline{R}) d\underline{R}' = 1 \quad (50)$$

$$\int_{V^N} \delta(\underline{R}' - \underline{R}) d\underline{R} = 1$$