

## 8. Mass, Radius and Luminosity Relations

### 8.1 The Eddington Luminosity

Assume radiation pressure.

$$\begin{aligned}
 P_{rad} &= \frac{1}{3} a T^4 \\
 \frac{dP_{rad}}{dP} &= \frac{\kappa \ell}{4\pi c G m(r)} \quad (187) \\
 P &= P_{gas} + P_{rad} \\
 \frac{dP}{dP} &= \frac{dP_{gas}}{dP} + \frac{dP_{rad}}{dP} = 1 \\
 \frac{dP_{rad}}{dP} &< 1
 \end{aligned}$$

$$\left( \begin{array}{l}
 \frac{dP_r}{dr} = \frac{dP_r}{dP} \frac{dP}{dr} \\
 \frac{dP_{rad}}{dr} = \frac{dP}{dT} \frac{dT}{dr} = \frac{4}{3} a T^3 \frac{dT}{dr} \\
 \frac{dT}{dr} = -\frac{3}{4ac} \frac{\kappa \rho}{T^3} \frac{\ell}{4\pi r^2}
 \end{array} \right)$$

(Note that all pressures increase inward, therefore all  $dP_n / dP$ 's are positive).

$$\kappa \ell < 4\pi c G m(r) \quad (188)$$

At the stellar surface,

$$L < \frac{4\pi c G M}{\kappa} \quad (189)$$

For opacity equal to electron scattering,

$$\kappa = \kappa_{es} = \text{const}$$

$$L < L_{Ed} = 3.2 \times 10^4 \left( \frac{M}{M_\odot} \right) L_\odot \quad (190)$$

Highest luminosity a stable star can have.

### 8.2 The Mass-Luminosity Relation

The scaling relations are implied by the equations of stellar structure. (The proper derivations use a technique called homology).

Hydrostatic equilibrium equation:

$$\frac{dP}{dr} = -g\rho = -\frac{GM(r)}{r^2} \rho \quad (191)$$

Radiation transport equation:

$$\frac{dT}{dr} = -\frac{\ell(r)}{4\pi r^2} \frac{3}{16} \frac{\kappa \rho}{\sigma T^3} \quad (192)$$

Approximate the density by its average,  $\bar{\rho}$ .

$$\bar{\rho} \propto \frac{M_*}{R_*^3}$$

Replace the derivative with the ratio  $P / R_*$

$$\frac{P}{R_*} = \frac{GM_*}{R_*^2} \frac{M_*}{R_*^3} \propto \frac{M_*^2}{R_*^5}$$

The ideal gas law,  $P = \rho kT / \mu$ , gives

$$T \propto \frac{P}{\bar{\rho}} \propto \frac{M_*^2}{R_*^4} \frac{R_*^3}{M_*} \propto \frac{M_*}{R_*}$$

Using equation (192),

$$L \propto -R_*^2 T^3 \frac{dT}{dr} \frac{1}{\kappa \rho}$$

Replace derivative with ratio,

$$\frac{dT}{dr} = -\frac{T}{R_*}$$

and we get

$$L \propto \frac{M_*^3}{R_*^3} \frac{R_*^3}{M_*} \frac{M_*}{R_*^2} R_*^2 \propto M_*^3 \quad (193)$$

### 8.3 Other scaling relations

$$\frac{dP}{dr} = -\frac{Gm(r)}{r^2} \rho$$

$$\frac{P}{R} \sim \frac{M}{R^2} \frac{M}{R^3} \quad (194)$$

$$\frac{dm}{dr} = 4\pi r^2 \rho$$

$$\frac{M}{R} \sim R^2 \rho \quad (195)$$

$$\frac{d\ell}{dr} = 4\pi r^2 \rho \epsilon$$

$$\frac{L}{R} \sim R^2 \frac{M}{R^3} \epsilon \quad (196)$$

$$\ell(r) = \frac{4\pi r^2}{\rho} \frac{c}{\kappa} \frac{d}{dr} \left( \frac{aT^4}{3} \right)$$

$$L \sim R^2 \frac{R^3}{M} \frac{1}{\kappa} \frac{T^4}{R} \quad (197)$$

where the luminosity equation assumes radiative energy transport.

We also have some supplementary equations,

$$P_g = \frac{nkT}{\mu}$$

$$P \sim \frac{M}{R^3} T$$

$$P_r = \frac{1}{3} a T^4$$

$$P \sim T^4$$

$$\varepsilon = \varepsilon_0 \rho T^n$$

$$\kappa = \kappa_0 \rho T^{-7/2}$$

As an example, if we assume  $\varepsilon = \varepsilon_0 \rho T^{15}$  and constant opacity (i.e. use equation 193),

$$M^{14} \propto R^{18} \quad (198)$$

$$R \propto M^{0.78} \quad (199)$$

For other values of n,

$$R \propto M^{(n-1)/(n+3)} \quad (200)$$

Use this to derive an  $(L, T_{eff})$  relation.

$$L \propto M^3$$

$$L = 4\pi R^2 T_{eff}^4$$

Note:  $T_{eff}$  is the surface temperature.  $T$  is the core temperature.

$$L^{1-(2/3)(n-2)/(n+3)} \propto T_{eff}^4$$

For  $n = 4$  (pp chain)

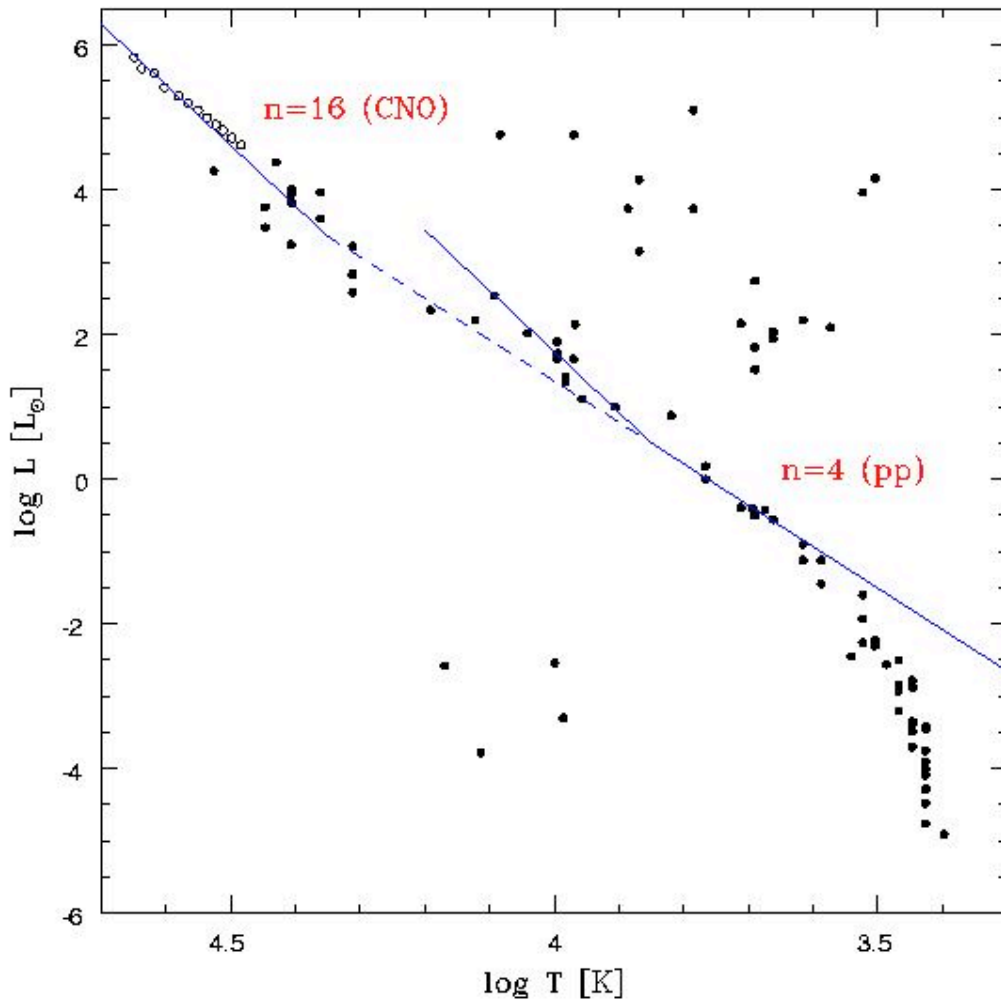
$$\log L = 5.6 \log T_{eff} + const \quad (201)$$

For  $n = 16$  (CNO chain)

$$\log L = 8.4 \log T_{eff} + const \quad (202)$$

For fully convective star ( $P \propto \rho^{4/3}$  and  $n = 4$ )

$$\log L = 3.7 \log T_{eff+const} \quad (203)$$



A reasonable fit is obtained to the observed Main Sequence.