

11. Time Independent Perturbation Theory

Objective: determine the eigenfunctions and eigenvalues of a Hamiltonian

$$\hat{H} = \hat{H}^0 + \lambda \hat{H}'$$

where \hat{H}^0 is a known solution for a similar problem, and \hat{H}' is the perturbation from this. The matrix elements of \hat{H}^0 and \hat{H}' are of the same size.

λ is a small number, so if \hat{H}^0 and \hat{H}' are the same size, the latter will give a small contribution to the result. λ is real, $|\lambda| \ll 1$.

We assume that we know the eigenfunctions and eigenvalues of \hat{H}^0 ,

$$\hat{H}^0 \varphi_n^0 = E_n^0 \varphi_n^0$$

with φ_n^0 orthonormal and E_n^0 non-degenerate.

We would like to solve

$$\hat{H} \varphi_n = E_n \varphi_n \quad (11-1)$$

As λ is a small number, we can expand around it.

$$\varphi_n = \varphi_n^0 + \lambda \varphi_n' + O(\lambda^2)$$

$$E_n = E_n^0 + \lambda E_n' + O(\lambda^2)$$

We chose the phase of φ_n such that

$$\langle \varphi_n, \varphi_n^0 \rangle = 0 \quad (11-2)$$

$$(\hat{H}^0 + \lambda \hat{H}')(\varphi_n^0 + \lambda \varphi_n') = (E_n^0 + \lambda E_n')(\varphi_n^0 + \lambda \varphi_n')$$

At order λ^0 :

$$\hat{H}' \varphi_n^0 + \hat{H}^0 \varphi_n' = E_n^0 \varphi_n' + E_n' \varphi_n^0 \quad (11-3)$$

We take the scalar product with φ_n^0 :

$$\langle \varphi_n^0, \hat{H}' \varphi_n^0 \rangle + \langle \varphi_n^0, \hat{H}^0 \varphi_n' \rangle = E_n^0 \underbrace{\langle \varphi_n^0, \varphi_n' \rangle}_0 + E_n' \underbrace{\langle \varphi_n^0, \varphi_n^0 \rangle}_1 \quad (11-4)$$

$$\langle \varphi_n^0, \hat{H}^0 \varphi_n' \rangle = \left\langle \underbrace{\hat{H}^0 \varphi_n^0}_{E_n^0 \varphi_n^0}, \varphi_n' \right\rangle = E_n^0 \underbrace{\langle \varphi_n^0, \varphi_n' \rangle}_0$$

$$E_n' = \langle \varphi_n^0, \hat{H}' \varphi_n^0 \rangle$$

So we can say that the total energy is

$$\begin{aligned} E_n &= E_n^0 + \lambda E_n' \\ &= E_n^0 + \lambda \langle \varphi_n^0, \hat{H}', \varphi_n^0 \rangle \end{aligned}$$

Going back to (11-3), and taking the scalar product with φ_i^0 ($i \neq n$):

$$\langle \varphi_i^0, \hat{H}', \varphi_n^0 \rangle + \langle \varphi_i^0, \hat{H}^0, \varphi_n' \rangle = E_n^0 \langle \varphi_i^0, \varphi_n' \rangle + E_n' \langle \varphi_i^0, \varphi_n^0 \rangle \quad (11-6)$$

We know that $\langle \varphi_i^0, \varphi_n^0 \rangle = 0$, and $\langle \varphi_i^0, \hat{H}^0, \varphi_n' \rangle = \langle \hat{H}^0 \varphi_i^0, \varphi_n' \rangle = E_i^0 \langle \varphi_i^0, \varphi_n' \rangle$. So,

(11-6) can be written as:

$$\langle \varphi_i^0, \hat{H}', \varphi_n^0 \rangle = E_n^0 \langle \varphi_i^0, \varphi_n' \rangle - E_i^0 \langle \varphi_i^0, \varphi_n' \rangle$$

$$\langle \varphi_i^0, \varphi_n^1 \rangle = \frac{\langle \varphi_i^0, \hat{H}^1, \varphi_n^0 \rangle}{E_n^0 - E_i^0} \quad (11-7)$$

φ_n^1 is an orthonormal basis of the state space. We can write any eigenfunction as a linear combination:

$$\varphi_n^1 = \sum_p c_p \varphi_p^0,$$

where

$$c_p = \langle \varphi_p^0, \varphi_n^1 \rangle.$$

We know that $\langle \varphi_n^0, \varphi_n^1 \rangle = 0$, hence $c_n = 0$. So we can write

$$\varphi_n^1 = \sum_{p \neq n} c_p \varphi_p^0$$

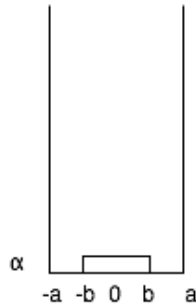
Using the equation for $\langle \varphi_i^0, \varphi_n^1 \rangle$ above, we can write:

$$\varphi_n^1 = \sum_{p \neq n} \frac{\langle \varphi_p^0, \hat{H}^1, \varphi_n^0 \rangle}{E_n^0 - E_p^0} \varphi_p^0$$

So we can write

$$\varphi_n = \varphi_n^0 + \lambda \varphi_n^1$$

Example: find the eigenfunctions and eigenvalues to the first order α of the Hamiltonian corresponding to the potential



$\hat{H}^0 =$ infinite square well. $|\alpha| \ll 1$

$$\hat{H} = \hat{H}^0 + \alpha \hat{H}^1$$

$$\hat{H}^1 = \begin{cases} 1 & |x| \leq b \\ 0 & |x| > b \end{cases}$$

We know that the energy for an infinite square well is:

$$E_n^0 = \frac{n^2 \pi^2 \hbar^2}{8a^2 m}$$

$$\varphi_n^0 = \begin{cases} \frac{1}{\sqrt{a}} \cos \frac{n\pi x}{2a} & n \text{ odd} \\ \frac{1}{\sqrt{a}} \sin \frac{n\pi x}{2a} & n \text{ even} \end{cases}$$

We said before that

$$E_n^1 = \langle \varphi_n^0, \hat{H}^1, \varphi_n^0 \rangle$$

Let's look at the correction to the energy level, with $n = 1$, and using

$\varphi_1^0 = \frac{1}{\sqrt{a}} \cos \frac{\pi x}{2a}$, we can write:

$$E_1' = \langle \varphi_1^0, \hat{H}' \varphi_1^0 \rangle = \frac{1}{a} \int_{-b}^b \cos \frac{\pi x}{2a} \left[1 - \frac{\cos \pi x}{2a} \right] dx = \frac{b}{a} + \frac{1}{\pi} \sin \frac{\pi b}{a}$$

So the first energy level is:

$$E_1 = \frac{\pi^2 \hbar^2}{8a^2 m} + \alpha \left(\frac{b}{a} + \frac{1}{\pi} \sin \left(\frac{\pi b}{a} \right) \right)$$

Now look at the second energy level.

$$E_2' = \langle \varphi_2^0, \hat{H}' \varphi_2^0 \rangle = \frac{1}{a} \int_{-b}^b \sin \frac{\pi x}{a} \sin \frac{\pi x}{a} dx = \frac{b}{a} - \frac{1}{2\pi} \sin \frac{2\pi b}{a}$$

So the second energy level is:

$$E_2 = \frac{4\pi^2 \hbar^2}{8a^2 m} + \alpha \left[\frac{b}{a} - \frac{1}{2\pi} \sin \frac{2\pi b}{a} \right]$$

Now look at $\langle \varphi_p^0, \hat{H}' \varphi_n^0 \rangle$. $p \neq n$.

$$\langle \varphi_p^0, \hat{H}' \varphi_n^0 \rangle = \int_{-b}^b \varphi_p^0 \varphi_n^0 dx$$

If n is odd, then φ_n will be even. If p is even, then φ_p is odd.

The scalar product will be 0 if p and n are of different parity (i.e. one odd, one even).

It will be non-zero if p and n are either both even, or both odd.

So, if n is odd, p will have to be odd, and φ_n' (from (11-7)) is a linear combination of even functions, so it is even. The same applies for odd functions.